

KUMMER RIGIDITY FOR K3 SURFACE AUTOMORPHISMS VIA RICCI-FLAT METRICS

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ABSTRACT. We give an alternative proof of a result of Cantat & Dupont, showing that any automorphism of a K3 surface with measure of maximal entropy in the Lebesgue class must be a Kummer example. Our method exploits the existence of Ricci-flat metrics on K3s and also covers the non-projective case.

1. INTRODUCTION

A basic result of Yomdin [Yom87], known previously as the Shub Entropy Conjecture, says that the topological entropy of any smooth map of a compact manifold is bounded below by the spectral radius of the action on homology. Gromov [Gro03] showed that in fact for compact Kähler manifolds and holomorphic automorphisms, this lower bound is always achieved. Thus entropy can be computed from linear-algebraic data, and when it is positive the measure of maximal entropy is unique on compact Kähler surfaces by [DTD12, Thm. 1.2]. For an introduction to complex dynamics in higher dimensions see [FS94], [Sib99], or [Can14].

Kummer examples are tori with automorphisms that become affine on the universal cover, together with their modifications using basic operations of birational geometry. A characteristic feature of Kummer examples is that the measure of maximal entropy is in the Lebesgue class. In [CD15], Cantat-Dupont proved that automorphisms with positive topological entropy of *projective* surfaces with measure of maximal entropy in the Lebesgue class are Kummer examples.

The main goal of this article is to give an alternative proof in the case of K3 surfaces, which in addition covers the non-projective case.

1.1. Main statements.

1.1.1. Theorem. *Let X be a K3 surface, $T: X \rightarrow X$ an automorphism with positive topological entropy whose measure of maximal entropy is absolutely continuous with respect to the Lebesgue measure. Then X is a Kummer K3 and T is induced by an (affine) automorphism of the corresponding torus.*

This resolves a conjecture of Cantat [Can99a, p.162] and McMullen [McM03, Conjecture 3.31], including the case of non-projective $K3$ surfaces. Combining [CD15] with our main theorem we easily obtain the following generalization:

1.1.2. Corollary. *Let X be a compact complex surface, $T: X \rightarrow X$ an automorphism with positive topological entropy whose measure of maximal entropy is absolutely continuous with respect to the Lebesgue measure. Then (X, T) is a Kummer example.*

As another application, our main result implies that the measure of maximal entropy in McMullen's construction of a Siegel disc [McM02] (which are never projective) cannot be in the Lebesgue class. Whether the complement of the Siegel disc can have positive Lebesgue measure remains an interesting open question.

For rational maps of \mathbb{P}^1 an analogous result was established by Zdunik [Zdu90], and for general endomorphisms of \mathbb{P}^n by Berteloot–Dupont [BD05] and Berteloot–Loeb [BL01], with the role of Kummer examples now played by Lattès maps. For related results in the case of general endomorphisms of Kähler manifolds see [Can08]. For an introduction to $K3$ surfaces, including their Ricci-flat metrics, see [K3-85].

1.2. Proof Outline. Let μ be the measure of maximal entropy and $d\text{Vol}$ the normalized volume form induced by the holomorphic 2-form on the $K3$ surface. The assumption says that $\mu = f d\text{Vol}$ for some $f \in L^1(d\text{Vol})$.

Because $d\text{Vol}$ is invariant under any holomorphic automorphism and since μ is ergodic (in fact mixing by [Can01, Thm. 6.1]), it follows that f is the normalized indicator function of a set of positive Lebesgue measure and hence we can assume

$$(1.2.1) \quad \mu = \frac{1}{d\text{Vol}(S)} d\text{Vol}|_S$$

for some T -invariant set S of positive Lebesgue measure. The proof is then naturally divided into three separate steps:

- *Step 1:* We prove the result in the special case $\mu = d\text{Vol}$ (i.e. $S = X$ above).
- *Step 2:* We prove the result when μ is uniformly hyperbolic and in the Lebesgue class, by reducing to the previous case.
- *Step 3:* We prove that μ is uniformly hyperbolic, assuming that it is in the Lebesgue class.

Step 1 is handled in Section 3 using the Ricci-flat Kähler metric provided by Yau's theorem [Yau78]. This key input, combined with Jensen's inequality, allows for an elementary proof. It already covers the case when the stable/unstable eigencurrents are smooth (since the indicator function of the set S has to be smooth, hence $S = X$), which was conjectured by Cantat [Can99a, p.162]. The two key properties that are used throughout

are that the metric is Kähler (so certain integrals are computed cohomologically) and that the metric induces the same volume form as the holomorphic 2-form. Each of these conditions can be easily ensured individually, but their simultaneous validity is Yau’s theorem.

Step 2 is handled in [Section 4](#) and is quite general: it would apply to smooth volume-preserving complex surface diffeomorphisms which are uniformly hyperbolic on an ergodic component of Lebesgue measure.

Finally, Step 3 is handled in [Section 5](#). We show that the expansion/contraction coefficients of the dynamics are in fact cohomologous to a constant (in the dynamical sense) and then show that the coboundaries are, in fact, uniformly bounded. We use that $\mu = \eta_+ \wedge \eta_-$ for two positive closed currents η_{\pm} called the stable/unstable eigencurrents. The coboundary structure follows from comparing the conditional measures of μ on the stable/unstable foliations and the restriction of η_{\pm} to these foliations. To prove the L^{∞} bound for the coboundaries, we first show that the restriction of the stable/unstable eigencurrent to a stable/unstable manifold is flat. This is then used to derive a contradiction if the uniform bound does not hold by constructing a Brody curve (in an orbifold K3 surface which is obtained by contracting some (-2) -curves in X) which intersects trivially both the stable and unstable eigencurrents. Such an entire curve is known not to exist by work of Dinh–Sibony [[DS05](#)]. This idea is also a key input in the work of Cantat–Dupont [[CD15](#)].

[Section 6](#) contains an alternative derivation of the coboundary structure, based again on Ricci-flat metrics, some estimates of Birkhoff sums, and ideas along the lines of the Gottschalk–Hedlund theorem. These arguments give *a priori* exponential integrability of the coboundaries, and it is possible that one could prove ergodicity of $d\text{Vol}$ using this weaker property, rather than the L^{∞} bound used in [Section 4](#).

Orbifolds. Because the automorphism potentially (in the Kummer case, always) has periodic curves, these must be contracted by a map $\nu: X \rightarrow Y$ and many arguments happen on Y instead of X . In the K3 case Y has only orbifold singularities and the needed properties are explained in [§2.1](#). On a first reading, one can assume that $X = Y$ and skip [Proposition 2.1.5](#).

Comparison to the approach of Cantat–Dupont. Our use of Ricci-flat metrics on K3 surfaces allows for several simplifications compared to the approach in [[CD15](#)]. A key step in showing that the K3 is Kummer is based on the existence of expanded/contracted foliations with the required smoothness properties. In our case, this follows directly from the equality case in Jensen’s inequality. Moreover, the Ricci-flat metric already provides the flat metric on the torus (see [§3.2](#)) and leads to an alternative proof that we have a Kummer example. Our arguments avoid the λ -lemma and holonomy maps that appear in [[CD15](#)].

To establish ergodicity (Step 2), we use a standard tool in dynamics, the Hopf argument (in a quantitative form). [[CD15](#)] is based on a topological

approach, showing that the set where the foliations are defined is both open and closed.

Finally, the use of Dinh–Sibony’s [Proposition 5.2.12](#) and the construction of a Brody curve on which the stable/unstable currents vanish is common to both proofs. We use it to establish uniform hyperbolicity of the measure of maximal entropy, while [\[CD15\]](#) use it to obtain a compactness property of the family of stable/unstable manifolds. The affine structure on these manifolds, a standard fact in dynamics, also appears in both proofs. The use of the contraction $\nu: X \rightarrow Y$ is unavoidable in both proofs.

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2. PRELIMINARIES

In this section we collect some preliminaries discussions, mostly well-known, and fix notation for the rest of the paper. On a first reading, the construction of the orbifold Y in [Proposition 2.1.5](#) and its Ricci-flat metrics can be skipped and one can assume that $X = Y$. An introduction to dynamics on surfaces can be found in the survey [\[Can14\]](#).

2.1. Preliminaries on currents and orbifolds.

2.1.1. *The eigencurrents.* For a detailed discussion of the results quoted in this section, see [\[Can01\]](#).

Let T be an automorphism with positive topological entropy $h > 0$ of a complex $K3$ surface X . Then there exist two closed positive $(1, 1)$ -currents η_{\pm} on X which satisfy

$$(2.1.2) \quad T^* \eta_{\pm} = e^{\pm h} \eta_{\pm}$$

and have continuous local potentials, in fact Hölder continuous by Dinh–Sibony [\[DS05\]](#). The currents yield cohomology classes $[\eta_{\pm}]$ and we normalize them to have $\int_X [\eta_+] \wedge [\eta_-] = 1$. These classes are nef, i.e. are limits of Kähler cohomology classes, and have vanishing self-intersection $\int_X [\eta_{\pm}]^2 = 0$.

We will be particularly interested in the cohomology $(1, 1)$ -class $[\eta_+] + [\eta_-]$ which is also clearly nef and satisfies $\int_X ([\eta_+] + [\eta_-])^2 = 2 \int_X [\eta_+] \wedge [\eta_-] = 2$.

2.1.3. *The picture in $H^{1,1}$.* The intersection pairing on $H^{1,1}(X, \mathbb{R})$ has signature $(1, 19)$ and the class $[\eta_+] + [\eta_-]$ sits on one of the two hyperboloids of classes with square 2. The hyperboloid is naturally identified with hyperbolic 19-space and contains a geodesic determined by intersecting with the two-dimensional real space spanned by $[\eta_+], [\eta_-]$. Cup product has signature $(1, 1)$ on this plane, and is negative definite on its orthogonal complement. The pullback action of the automorphism T acts as an isometry on the complement to the 2-plane, and as a translation by h along the geodesic that the plane determines.

The Kähler classes in $H^{1,1}(X, \mathbb{R})$ are those that pair positively against (-2) curves (see [K3-85, XIII.2]) and the Kähler cone contains the geodesic in its closure. Typically the geodesic is strictly in the interior and the orbifold construction below is unnecessary (this can be assumed on a first reading). When the geodesic lies in the boundary, one has to contract the (-2) curves that pair to zero against both $[\eta_+], [\eta_-]$ as in Proposition 2.1.5 below.

For more on constructing automorphisms of K3 surfaces using this point of view, see [McM11].

2.1.4. *The associated orbifold.* The following result is well-known, but for the reader's convenience we provide a proof.

2.1.5. **Proposition.** *Let $V \subset X$ be the union of all irreducible compact holomorphic curves $C \subset X$ which satisfy $\int_C([\eta_+] + [\eta_-]) = 0$.*

- (i) *There are finitely many compact holomorphic curves periodic¹ under T , and V is their union.*
- (ii) *There exists an orbifold Y and a holomorphic map $\nu: X \rightarrow Y$, which is an isomorphism away from V and contracts each connected component of V to an orbifold point of Y .*
- (iii) *There exists a holomorphic automorphism T_Y of Y with the same topological entropy as T , such that $\nu \circ T = T_Y \circ \nu$.*

From this result it follows that $X = Y$ if and only if T has no periodic curves, or equivalently the class $[\eta_+] + [\eta_-]$ is Kähler on X .

Proof. For part (i), it is clear that T -periodic curves are contained in V , since $\int_C \eta_{\pm} = \int_{T^{-N}C} (T^N)^* \eta_{\pm} = e^{\pm Nh} \int_{T^{-N}C} \eta_{\pm}$ for any compact curve C and integer N . A similar calculation shows that V is T -invariant and the general result in [CT15, Theorem 1.1] implies that V is the union of finitely many curves, proving (i).

For part (ii), since the intersection form is negative definite on the complement of the span of $[\eta_+], [\eta_-]$, it follows that if we write $V = \cup_{i=1}^N C_i$ for the decomposition of V into irreducible components, then the intersection matrix $(C_i \cdot C_j)$ is negative definite. By a theorem of Grauert [Gra62] there

¹A T -periodic curve is defined to satisfy $T(C) = C$ set-theoretically, but not necessarily pointwise.

is then a contraction map $\nu: X \rightarrow Y$ onto an irreducible normal compact complex surface, which contracts each connected component of V to a point.

Next we claim that each connected component of V is in fact an *ADE* curve (i.e. its irreducible components are smooth rational curves with self-intersection equal to -2 [BHPV04, §III.2]). Indeed, the adjunction formula gives

$$p_a(C_i) = 1 + \frac{(K_X \cdot C_i) + (C_i^2)}{2} \leq \frac{1}{2},$$

using of course that $K_X \cong \mathcal{O}_X$, hence $p_a(C_i) = 0$ and so each C_i is a smooth rational curve with $(C_i^2) = -2$, as claimed.

The fact that V is composed of *ADE* curves now implies that all singular points of Y are rational double points, which in particular are orbifold points (locally isomorphic to the quotient \mathbb{C}^2/Γ for certain finite subgroups $\Gamma \subset \mathrm{SU}(2)$ acting freely on the unit sphere, see [Dur79] for more).

For part (iii), following e.g. [Kaw08, Lemma 2.2] since T maps V onto itself, it descends to an automorphism $T_Y: Y \rightarrow Y$ with $\nu \circ T = T_Y \circ \nu$, and the topological entropy of T_Y equals the one of T , namely h . \square

2.1.6. The invariant measures. Because the currents η_{\pm} have continuous local potentials, their wedge product in the sense of Bedford-Taylor [BT76] is well-defined and gives a T -invariant probability measure $\mu := \eta_+ \wedge \eta_-$ on X . It is the unique measure of maximal entropy.

On the other hand, the K3 surface X carries a nowhere vanishing holomorphic 2-form Ω , which we normalize to have $\int_X \Omega \wedge \bar{\Omega} = 1$. The probability measure $\mathrm{dVol} := \Omega \wedge \bar{\Omega}$ is automatically T -invariant, and will be referred to as the Lebesgue measure of X . It will sometimes be denoted by dVol_X .

Applying ν_* gives a T_Y -invariant volume form $\mathrm{dVol}_Y := \nu_*(\mathrm{dVol}_X)$, which is nowhere vanishing in the orbifold sense.

2.1.7. Orbifold Ricci-flat metrics. The orbifold K3 surface Y constructed in Proposition 2.1.5 admits orbifold Kähler metrics, as is well-known (see e.g. [Cam04, Ex. 3.2]). In fact, for every $t \in \mathbb{R}$ there is an orbifold Kähler class $[\omega_{Y,t}]$ on Y such that

$$(2.1.8) \quad e^t[\eta_+] + e^{-t}[\eta_-] = \nu^*[\omega_{Y,t}]$$

holds. The (nontrivial) proof of this fact is given in [FT18, Proof of Theorem 1.3]. In our setting, Ricci-flatness of a (normalized) Kähler metric ω_Y is equivalent to $\omega_Y^2 = \mathrm{dVol}_Y$.

2.1.9. Proposition. *With the notation as in Eqn. (2.1.8), we have:*

- (i) *The Kähler class $[\omega_{Y,t}]$ on the orbifold Y contains a unique orbifold Ricci-flat Kähler metric $\omega_{Y,t}$.*
- (ii) *The metrics satisfy $\omega_{Y,t+h} := T_Y^*(\omega_{Y,t})$.*

Proof. For (i), the proof of Yau's theorem [Yau78] extends to orbifolds (see e.g. [Cam04, Thm. 4.1]), and so we conclude that the class $[\omega_{Y,t}]$ contains a unique orbifold Ricci-flat Kähler metric $\omega_{Y,t}$.

Part (ii) follows by uniqueness of the Ricci-flat metric in its cohomology class: both $\omega_{Y,t+h}$ and $T_Y^*(\omega_{Y,t})$ are Ricci-flat and we have

$$\begin{aligned} \nu^*[T_Y^*(\omega_{Y,t})] &= T^*\nu^*[\omega_{Y,t}] = T^*(e^t[\eta_+] + e^{-t}[\eta_-]) = e^{t+h}[\eta_+] + e^{-t-h}[\eta_-] \\ &= \nu^*[\omega_{Y,t+h}], \end{aligned}$$

and the map ν^* is injective in cohomology. \square

2.2. General remarks about Lyapunov exponents. Let Z denote either X or Y , and T_Z the corresponding automorphism. Let m be any T_Z -invariant probability measure. Fix a smooth hermitian metric for computing all norms below and recall that the Lyapunov exponent of m is defined by (see [Led84, §1.1]):

$$\lambda(m) := \lim_{N \rightarrow \infty} \frac{1}{N} I_N \text{ with } I_N := \int_Z \log \|DT_Z^N\| dm$$

The limit exists since I_n satisfies the subadditivity property $I_{k+l} \leq I_k + I_l$:

$$\begin{aligned} \int_Z \log \|D_x T^{k+l}\| dm(x) &\leq \int_Z \log \left(\|D_x T^k\| \cdot \|D_{T^k x} T^l\| \right) dm(x) \\ &= \int_Z \log \|D_x T^k\| dm(x) + \int_Z \log \|D_{T^k x} T^l\| dm(x) \\ &= I_k + \int_Z \log \|D_x T^l\| dm(T^{-k}x) = I_k + I_l \end{aligned}$$

where we have used the T_Z -invariance of m and the inequality $\|A \cdot B\| \leq \|A\| \cdot \|B\|$ for linear maps A, B . In fact by Fekete's lemma $\lambda(m) = \inf_N \frac{1}{N} I_N$ and in particular

$$(2.2.1) \quad \frac{1}{N} I_N \geq \lambda(m) \quad \forall N \geq 1.$$

The exponent does not depend on the fixed ambient metric, since any two will be uniformly comparable.

2.2.2. Stable/Unstable directions. If m is a T_Z -invariant ergodic probability measure with strictly positive Lyapunov exponent, then the Oseledets theorem (see e.g. [Led84, Thm. 3.1]) implies that there exist measurable DT_Z -invariant complex line subbundles $W^\pm(x)$ of the tangent bundle of Y , defined for m -a.e. x , such that $\lim_{N \rightarrow \infty} \frac{1}{N} \log \|D_x T_Z^N|_{W^\pm(x)}\| = \pm \lambda(m)$ for m -a.e. x (the positive and negative exponents have the same absolute value because T_Z is volume-preserving). We will use alternatively the notation W^s for W^- and call it the stable direction, and W^u for W^+ for the unstable.

2.2.3. Absolute continuity. Suppose now that the measure of maximal entropy μ is absolutely continuous with respect to the Lebesgue measure. Then Ledrappier–Young [LY85, Corollary G] implies that the Lyapunov exponent of μ is $\frac{h}{2}$, since the real dimension of the unstable subspace is 2 and the entropy is h .

3. THE ERGODIC CASE

In this section we assume that in fact $\mu = \text{dVol}$ and give an easy proof that (X, T) is a Kummer example, using Ricci-flat metrics and the equality case of Jensen's inequality.

3.1. Using Jensen's inequality. For the orbifold Ricci-flat metrics on Y constructed in Proposition 2.1.9, let $\omega_Y := \omega_{Y,0}$ be the fixed reference metric.

3.1.1. Definition (Expansion factor). For $x \in Y$ let $\lambda(x, N)$ denote the expansion factor (or pointwise Lyapunov exponent) after N iterates of the map T_Y . Namely, at a given point x the metrics ω_Y and $\omega_{Y, Nh}$ are of the form, in an appropriate basis:

$$(3.1.2) \quad \begin{aligned} \omega_Y(x) &= |dz_1|^2 + |dz_2|^2 \\ \omega_{Y, Nh}(x) &= |e^{\lambda(x, N)} dz_1|^2 + |e^{-\lambda(x, N)} dz_2|^2 \end{aligned}$$

More intrinsically, we can define $\lambda(x, N)$ to be equal to $\frac{1}{2} \log$ of the largest eigenvalue of the hermitian form $\omega_{Y, Nh}(x)$ with respect to $\omega_Y(x)$.

The following simple observation is the key which yields the main theorem in the case when $\mu = \text{dVol}$. Note that the assumption is on X , but the conclusions are on Y .

3.1.3. Proposition. *Assume that $\mu = \text{dVol}$ on X . Then*

- (i) *We have $\lambda(x, N) = \frac{Nh}{2}$ for every $x \in Y$ and every N .*
- (ii) *There exist at every $x \in Y$ two orthogonal (for ω_Y) tangent directions $W^\pm(x)$ such that $\log \|DT_Y|_{W^\pm(x)}\| = \pm \frac{h}{2}$. The directions $W^\pm(x)$ vary real-analytically in x , are T_Y -invariant, and agree with the directions provided by the Oseledets theorem applied to μ .*

Proof. For part (i), the cohomological calculation $\int_Y \omega_Y \wedge \omega_{Y, Nh} = e^{Nh} + e^{-Nh}$ and Jensen's inequality give:

$$(3.1.4) \quad \begin{aligned} \log(e^{Nh} + e^{-Nh}) &= \log \left(\int_Y \omega_Y \wedge \omega_{Y, Nh} \right) \\ &\geq \int_Y \log \left(\frac{\omega_Y \wedge \omega_{Y, Nh}}{\text{dVol}_Y} \right) \text{dVol}_Y \\ &= \int_Y \log(e^{2\lambda(x, N)} + e^{-2\lambda(x, N)}) \text{dVol}_Y. \end{aligned}$$

We established in Eqn. (2.2.1) that

$$\int_Y \lambda(x, N) \text{dVol}_Y \geq \frac{Nh}{2}.$$

Indeed, in the case at hand $\log \|D_x T_Y^N\| = \lambda(x, N)$ by definition, and the Lyapunov exponent is $h/2$ by §2.2.3.

Note that the function $\log(e^x + e^{-x})$ is convex and increasing so we can apply Jensen again. Recall that $2I_N := \int_Y 2\lambda(x, N) \, d\text{Vol}_Y \geq Nh$ to find:

$$(3.1.5) \quad \int_Y \log(e^{2\lambda(x, N)} + e^{-2\lambda(x, N)}) \, d\text{Vol}_Y \geq \log(e^{2I_N} + e^{-2I_N}) \geq \log(e^{Nh} + e^{-Nh})$$

So from Eqn. (3.1.4) and Eqn. (3.1.5) it follows that we must have equality pointwise a.e., that is $\lambda(x, N) = Nh/2$ pointwise a.e., for all N . Since the function $\lambda(-, N)$ is continuous, the result holds everywhere on Y .

For part (ii), the equality case in Jensen plus the equality case in $I_k + I_\ell \leq I_{k+\ell}$ (see §2.2) imply that at every point $x \in Y$, the directions dz_1, dz_2 appearing in Eqn. (3.1.2) are independent of N and determine the spaces $W^\pm(x)$; they are orthogonal for ω_Y by the spectral theorem for hermitian matrices. Moreover, for any real-analytic Kähler metrics α, β , the direction of maximal expansion of α relative to β varies real-analytically, away from the locus where the direction is not unique (empty in our case). \square

We can now finish off the proof of Theorem 1.1.1 when $\mu = \text{dVol}$.

3.1.6. Theorem. *Let X be a K3 surface, $T: X \rightarrow X$ an automorphism with positive topological entropy whose measure of maximal entropy μ equals Lebesgue measure dVol . Then X is a Kummer K3 and T is induced by an (affine) automorphism of the corresponding torus.*

Proof. Proposition 3.1.3(ii) gives two line subbundles of the tangent bundle of Y , invariant and uniformly expanded/contracted by the dynamics. By Ghys [Ghy95, Proposition 2.2] [CD15, Proposition 7.1], these give two holomorphic foliations on the orbifold Y . These can be pulled back to holomorphic foliations on $X \setminus V$, which automatically extend to X exactly as in [CD15, Proof of Corollary 7.6]), which are preserved by T (alternatively, as Cantat pointed out to us, in the present case the extension of the holomorphic foliations follows from the explicit description of the singular points of Y as quotient ADE singularities). At this point we can apply either a result of Cantat [Can01, Theorem 7.4], or a later result of Cantat-Favre [CF03, Theorem 3.1] (which only needs one invariant foliation), to conclude that X is a Kummer K3 and T is induced by an automorphism of the corresponding torus. Note that these results apply in the Kähler case (for [CF03] one needs to use results of Brunella [Bru06, Bru07]). \square

3.2. An alternative argument. At the end of the proof of Theorem 3.1.6 above, after obtaining two T -invariant holomorphic foliations on X , we appealed to the general results of Cantat [Can01] or Cantat-Favre [CF03] to conclude that (X, T) is a Kummer example.

We now explain how to circumvent in our case some of the just cited arguments using the differential geometry of Ricci-flat metrics. We suppose that on X we have $\mu = \text{dVol}$, and we have applied Proposition 3.1.3, to obtain two T_Y -invariant transverse holomorphic foliations on Y , as above.

3.2.1. Proposition. *In this setting, the orbifold Ricci-flat Kähler metric ω_Y on Y is in fact flat.*

Proof. At any point $x \in Y$ we have the two orbifold Ricci-flat Kähler metrics ω_Y and $\omega_{Y,h} = T_Y^* \omega_Y$, which have the property that the eigenvalues of $\omega_{Y,h}$ with respect to ω_Y are e^h and e^{-h} , and the corresponding eigenvectors (these are eigenvectors of the endomorphism of the tangent space given by composing $\omega_{Y,h}$ with ω_Y^{-1}) span the stable and unstable holomorphic foliations \mathcal{F}_\pm respectively. By construction, these eigenvectors are ω_Y -orthogonal, hence so are \mathcal{F}_+ and \mathcal{F}_- . Near x we can find local holomorphic functions z_1, z_2 such that $\ker dz_1 = \mathcal{F}_+$, $\ker dz_2 = \mathcal{F}_-$, which implies that z_1, z_2 give local holomorphic coordinates near x (on the local orbifold cover if x is singular). Therefore near x we can write

$$\omega_Y = a idz_1 \wedge d\bar{z}_1 + b idz_2 \wedge d\bar{z}_2,$$

where a, b are local smooth positive functions. But then

$$0 = d\omega_Y = \left(\frac{\partial a}{\partial z_2} dz_2 + \frac{\partial a}{\partial \bar{z}_2} d\bar{z}_2 \right) \wedge idz_1 \wedge d\bar{z}_1 + \left(\frac{\partial b}{\partial z_1} dz_1 + \frac{\partial b}{\partial \bar{z}_1} d\bar{z}_1 \right) \wedge idz_2 \wedge d\bar{z}_2,$$

which imply that a is independent of the z_2, \bar{z}_2 variables and b is independent of the z_1, \bar{z}_1 directions.

From the definition of curvature we have

$$R_{i\bar{j}k\bar{\ell}} = -\partial_i \partial_{\bar{j}} g_{k\bar{\ell}} + g^{p\bar{q}} \partial_i g_{k\bar{q}} \partial_{\bar{j}} g_{p\bar{\ell}},$$

where in our coordinates

$$g_{1\bar{2}} = g_{2\bar{1}} = g^{1\bar{2}} = g^{2\bar{1}} = 0, \quad g_{1\bar{1}} = a, g_{2\bar{2}} = b, \quad g^{1\bar{1}} = a^{-1}, g^{2\bar{2}} = b^{-1},$$

$$\partial_2 g_{1\bar{1}} = \partial_{\bar{2}} g_{1\bar{1}} = 0, \quad \partial_1 g_{2\bar{2}} = \partial_{\bar{1}} g_{2\bar{2}} = 0.$$

In particular

$$R_{2\bar{2}1\bar{1}} = -\partial_2 \partial_{\bar{2}} a + g^{p\bar{q}} \partial_2 g_{1\bar{q}} \partial_{\bar{2}} g_{p\bar{1}} = g^{1\bar{1}} \partial_2 g_{1\bar{1}} \partial_{\bar{2}} g_{1\bar{1}} = 0,$$

while Ricci-flatness gives

$$0 = R_{1\bar{1}} = g^{p\bar{q}} R_{p\bar{q}1\bar{1}} = g^{1\bar{1}} R_{1\bar{1}1\bar{1}} + g^{2\bar{2}} R_{2\bar{2}1\bar{1}} = a^{-1} R_{1\bar{1}1\bar{1}},$$

giving $R_{1\bar{1}1\bar{1}} = 0$, and

$$0 = R_{2\bar{2}} = g^{p\bar{q}} R_{p\bar{q}2\bar{2}} = g^{1\bar{1}} R_{1\bar{1}2\bar{2}} + g^{2\bar{2}} R_{2\bar{2}2\bar{2}} = b^{-1} R_{2\bar{2}2\bar{2}},$$

giving $R_{2\bar{2}2\bar{2}} = 0$ (using the Kähler identities). Next,

$$R_{1\bar{1}1\bar{2}} = -\partial_1 \partial_{\bar{1}} g_{1\bar{2}} + g^{p\bar{q}} \partial_1 g_{1\bar{q}} \partial_{\bar{1}} g_{p\bar{2}} = 0,$$

$$R_{1\bar{2}2\bar{2}} = -\partial_1 \partial_{\bar{2}} g_{2\bar{2}} + g^{p\bar{q}} \partial_1 g_{2\bar{q}} \partial_{\bar{2}} g_{p\bar{2}} = g^{2\bar{2}} \partial_1 g_{2\bar{2}} \partial_{\bar{2}} g_{2\bar{2}} = 0.$$

Thanks to the Kähler identities, we thus obtain that $R_{i\bar{j}k\bar{\ell}} = 0$ for all i, j, k, ℓ , hence ω_Y is a flat orbifold Kähler metric. \square

Once we know that ω_Y is flat, this in turn implies that there is a finite orbifold cover $\pi: Z \rightarrow Y$ with Z a compact complex 2-torus (cf. the discussion in [CZ12, §6.4.3]). The arguments in [Can01, CF03] can then be used to show that T_Y lifts to an automorphism of Z , which is then affine linear, and that the map $\pi: Z \rightarrow Y$ is the quotient by an involution, and so X is a Kummer K3 and T is induced by an automorphism of the corresponding torus.

4. THE UNIFORMLY HYPERBOLIC CASE

After recalling in §4.1 the needed facts from Pesin theory, we prove in §4.2 that if μ is uniformly hyperbolic and in the Lebesgue class, then it is in fact equal to $d\text{Vol}$.

4.1. Recollections from Pesin theory. This section collects some concepts and results from Pesin theory that will be used to prove ergodicity of $d\text{Vol}$, under an extra assumption of uniform hyperbolicity. While Pesin theory is concerned with the non-uniformly hyperbolic setting, its conclusions apply to sets of almost full measure, on which hyperbolicity is uniform.

For the discussion in this subsection, Z can be any compact complex surface, possibly an orbifold, $T_Z: Z \rightarrow Z$ a holomorphic automorphism, and m a T_Z -invariant ergodic probability measure with nonzero Lyapunov exponents in the sense of §2.2.

4.1.1. Definition (Uniform hyperbolicity). We say that m is *uniformly hyperbolic* if there exists a constant $C > 0$ such that for m -a.e. x , we have for $N \geq 1$ that

$$(4.1.2) \quad \log \left\| D_x T_Z^{-N} |_{W^u(x)} \right\| \leq -\frac{1}{C}N + C$$

$$(4.1.3) \quad \log \left\| D_x T_Z^N |_{W^s(x)} \right\| \leq -\frac{1}{C}N + C$$

and the angle between $W^u(x)$ and $W^s(x)$ is bounded below by $\frac{1}{C}$, where everything is measured relative to a fixed smooth Riemannian metric.

In [BP07, Def. 2.2.6] sets $\Lambda_{\lambda\mu\varepsilon j}^\ell$ are defined to which most considerations in loc.cit. apply. A T_Z -invariant set of points x which satisfy Definition 4.1.1 is then contained in such a $\Lambda_{\lambda\mu\varepsilon j}^\ell$ with $j = 1$, $\varepsilon = 0$ and ℓ, λ, μ only depending on C .

4.1.4. Stable and Unstable manifolds. The following discussion is expanded in [BP07, §7-8]. Because the invariant measure m has nonzero Lyapunov exponents, there exist for m -a.e. x unique global immersed stable manifolds $\mathcal{W}^s(x)$ which contain x and with tangent space equal to $W^s(x)$ there. For distinct points, stable manifolds either coincide or are disjoint.

Since in our case the stable manifolds are complex 1-dimensional they are parametrized by \mathbb{C} [BLS93, Prop. 2.6]. For convenience of notation, we

will use the canonical parametrizations given by the complex line $W^u(x)$, namely the holomorphic maps

$$\xi_x^s : W^s(x) \xrightarrow{\sim} \mathcal{W}^s(x)$$

normalized to have derivative the identity at the basepoints (where $0 \mapsto x$). Because the only holomorphic automorphisms of \mathbb{C} are affine, we immediately deduce that

- For two points x, y on the same stable manifold the composed map $(\xi_y^s)^{-1} \circ \xi_x^s$ is an affine map from $W^s(x)$ to $W^s(y)$.
- The maps ξ_x^s and $T_Z^{-1} \circ \xi_{T_Z(x)}^s \circ D_x T_Z$ coincide, because they both induce parametrizations of the stable manifolds at x and have the same derivative at the origin. Equivalently $T_Z \circ \xi_x^s = \xi_{T_Z(x)}^s \circ D_x T_Z$.

Therefore the stable manifolds carry canonical affine structures and are parametrized equivariantly for the dynamics. The same discussion applies to unstable manifolds when using T_Z^{-1} instead of T_Z , and we will denote their parametrizations by ξ_x^u .

4.1.5. *Charts and size.* To discuss the geometry of stable manifolds, fix finitely many open charts covering $Z = \cup_\alpha U_\alpha$ and view each chart as equipped with its flat Euclidean metric (in the orbifold sense when necessary). There exists an $\varepsilon > 0$ which is a Lebesgue number of this covering, i.e. for any $x \in Z$ there exists a chart U_α such that the ball of radius ε around Z is contained in U_α . All considerations below will be in these charts and all objects will be considered only in balls of radius at most $\varepsilon/2$, so that the Euclidean and fixed background metric are comparable, up to uniform constants.

Implicit constants occurring below will be called *uniform* if they only depend on the automorphism T_Z , the covering fixed above, and a fixed smooth ambient metric.

4.1.6. *Local stable manifolds.* Recall that the global stable manifolds $\mathcal{W}^s(x)$ are only immersed and are patched from local stable manifolds $\mathcal{W}_{loc}^s(x)$. The geometry of the local stable manifolds is described in [BP07, Thm. 7.5.1]. Most importantly, under the uniform hyperbolicity assumptions in Definition 4.1.1 the constants (that appear in loc.cit.) r giving the radius, and D giving the Hölder constant of the derivative, depend only on the constant describing the uniform hyperbolicity (by Thm. 7.5.1(5) of loc.cit.). As a consequence, using the charts from §4.1.5 to map a ball of radius r in $W^{s/u}(x)$ to Z , the stable resp. unstable manifolds will be contained in disjoint cones around $W^s(x)$ resp. $W^u(x)$, with angle between the cones uniformly bounded below.

4.2. Ergodicity in the uniformly hyperbolic case.

4.2.1. **Theorem.** *Suppose that the measure of maximal entropy μ on X is in the Lebesgue class and additionally under the contraction $\nu: X \rightarrow Y$*

from §2.1, the measure $\nu_*\mu$ is uniformly hyperbolic on Y in the sense of Definition 4.1.1.

Then $\mu = \text{dVol}$ on X , and so (X, T) is a Kummer example by Theorem 3.1.6.

The statements $\mu = \text{dVol}_X$ and $\nu_*\mu = \text{dVol}_Y$ are equivalent, since the set contracted by ν has Lebesgue measure zero. However, even when (X, T) is a Kummer example, dVol_X is not uniformly hyperbolic for a smooth metric on X .

Proof. Let $S \subset X$ be the set defined in Eqn. (1.2.1), so that $\mu = \frac{\text{dVol}}{|S|} \Big|_S$. All arguments below are on Y so for simplicity of notation let S denote the image of this set under ν .

Assume that the uniform hyperbolicity condition holds for every point of S (otherwise replace S with the intersection of all iterates of the set on which uniform hyperbolicity holds, still a set of full μ -measure). It suffices to show that there exists a uniform $\varepsilon > 0$ such that for Lebesgue-a.e. $x \in S$ the ball $B(x, \varepsilon)$ of radius ε around x satisfies $|B(x, \varepsilon) \cap S| = |B(x, \varepsilon)|$, i.e. the ball is essentially (up to Lebesgue measure 0) contained in S (we'll use $|B|$ to denote the volume of B for the canonical volume form).

This last property, in turn, follows from the Hopf argument and the uniformity of hyperbolicity. Specifically, from Pesin theory there exist stable and unstable manifolds $\mathcal{W}^s(x), \mathcal{W}^u(x)$ through μ -a.e. point x . From our uniform hyperbolicity assumption, the sizes of the stable/unstable manifolds in the sense of [BP07, §8.1] are uniformly bounded below for every $x \in S$, and so are their angles (this is [BP07, §8.1.3]). Indeed there exists a single $\ell > 0$ such that the Oseledets-regular level sets Λ_\bullet^ℓ that appear in [BP07] are equal to S . Furthermore, the stable and unstable manifolds depend continuously in the C^1 -topology for points in S , by [BP07, §8.1.4] (even better, Hölder continuity holds along the lines in [ABF16]).

Now we apply the Hopf argument. From [BP07, Lemma 9.1.2], there exists a set B_S of zero Lebesgue measure such that if $y \in S$ and $z, w \in \mathcal{W}^u(y) \setminus B_S$ then $z, w \in S$. Moreover, by the absolute continuity of the unstable foliations [BP07, Thm. 8.6.8], for μ -a.e. $y \in S$ the set $\mathcal{W}^u(y) \setminus B_S$ has zero Lebesgue measure in $\mathcal{W}^u(y)$. Let $S' \subset S$ be this last set, so that $|S \setminus S'| = 0$.

For $y \in S'$ we have that Lebesgue-a.e. $z \in \mathcal{W}^u(y)$ is in S , hence admits stable manifolds $\mathcal{W}^s(z)$ of size bounded below and which depend continuously on z . Hence we get a continuous injective map $\Delta \times \Delta \rightarrow Y$, where $\Delta = \{|z| < 1 : z \in \mathbb{C}\}$, as follows. The map is defined by $(z, w) \mapsto \xi_z^s(w)$ where the first factor of Δ is identified with a disc of size uniformly bounded below in $\mathcal{W}^u(y)$, and the second factor with a disc in $W^s(z)$ as z varies in $\mathcal{W}^u(y)$ (the disc Δ of radius 1 is rescaled from the bounded below radius r discussed in §4.1.6). While the map is defined only for Lebesgue-a.e. z , since

it is continuous as a map $z \mapsto \text{Hol}(\Delta, Y)$ (by [BP07, §8.1.4]), it extends to the desired continuous map $\Delta \times \Delta \rightarrow Y$.

By the Invariance of Domain Theorem ([Hat02, Thm. 2B.3]), this continuous injective map contains an open set around $y \in Y$. It remains to check that this open set contains a ball of size uniformly bounded below. For this, it suffices to check that there exists a uniform $\varepsilon > 0$ such that if $(z, w) \in \Delta \times \Delta$, $y' = \xi_z^s(w)$ and $|y - y'| \leq \varepsilon$ then $|z| + |w| \leq \frac{1}{10}$. Provided this last property, it is clear that every point within ε of y is in the image, since the set in question is both open (by invariance of domain) and closed by the property that remains to be checked.

Suppose therefore that $y' = \xi_z^s(w)$. Recall that in Euclidean triangles, by the law of sines if a is the length of one side and α is the opposite angle, then $\frac{a}{\sin \alpha}$ controls any other side. In the charts described in §4.1.5, the stable and unstable manifolds at z are contained in cones around $W^s(z)$ and $W^u(z)$ with angle between them uniformly bounded below (see §4.1.6). Therefore there exists a uniform constant A such that $A|y - y'| \geq |z| + |w|$, which suffices for our purposes. \square

5. THE GENERAL CASE

We complete the proof of the main theorem by establishing Step 3 from the introduction (§1.2).

5.1. Proving uniform hyperbolicity. Our goal is the following, which thanks to Theorem 3.1.6 (Step 1) and Theorem 4.2.1 (Step 2) completes the proof of Theorem 1.1.1:

5.1.1. Theorem. *Suppose that the measure of maximal entropy μ on X is in the Lebesgue class. Then under the contraction $\nu: X \rightarrow Y$, the measure $\nu_*\mu$ is uniformly hyperbolic on Y .*

Because sets contracted under $\nu: X \rightarrow Y$ have zero Lebesgue measure, μ -measurable functions on X are naturally identified with $\nu_*\mu$ -measurable ones on Y . We will not distinguish in notation between functions identified in this manner.

5.1.2. Expansion/Contraction factors. Recall (see §2.2.2 and §2.2.3) that $\nu_*\mu$ -a.e. $x \in Y$ has a decomposition of the tangent space:

$$\mathbf{T}_x Y = W^s(x) \oplus W^u(x)$$

which is T_Y -invariant. i.e. $DT_Y(W^{u/s}(x)) = W^{u/s}(T_Y x)$. Fix now the Calabi–Yau metric ω_Y on Y (see §2.1.7) and define the $\nu_*\mu$ -measurable functions:

$$(5.1.3) \quad \rho^u(x) := \log \|DT_Y|_{W^u(x)}\|_{\omega_Y}, \quad \rho^s(x) := \log \|DT_Y|_{W^s(x)}\|_{\omega_Y}$$

where the norm of the operator is for the metrics $\omega_Y(x)$ and $\omega_Y(T_Y x)$ on the source and target tangent spaces. The functions are $\nu_*\mu$ -measurable since the spaces $W^{u/s}$ depend measurably on x .

5.1.4. *Computation in a basis.* Let now $e_1 \in W^u(x), e_2 \in W^s(x)$ be a unimodular basis, i.e. using the dual basis $\sqrt{-1}e_1^\vee \wedge \overline{e_1}^\vee \wedge \sqrt{-1}e_2^\vee \wedge \overline{e_2}^\vee$ is the fixed volume form $d\text{Vol}_Y$. There are two metrics on $\mathbf{T}_x Y$, one is $\omega_Y(x)$ and the other one is $\omega_{Y, Nh}(x)$, which is the pull-back by T_Y^N of $\omega_Y(T_Y^N x)$.

Suppose that in the fixed basis $\{e_1, e_2\}$ the metrics are represented by

$$(5.1.5) \quad \omega_Y(x) = \begin{bmatrix} a_0 & b_0 \\ b_0 & d_0 \end{bmatrix}, \quad \omega_{Y, Nh}(x) = \begin{bmatrix} a_{Nh} & b_{Nh} \\ b_{Nh} & d_{Nh} \end{bmatrix}$$

and these matrices have determinant equal to 1. The relation to the functions defined in Eqn. (5.1.3) is:

$$\rho^u(x) = \frac{1}{2} \log \left(\frac{a_h}{a_0} \right), \quad \rho^s(x) = \frac{1}{2} \log \left(\frac{d_h}{d_0} \right)$$

5.1.6. *Defining β .* Consider now the quantity

$$(5.1.7) \quad \beta(x) := \frac{1}{2} \log(a_0 d_0)$$

where a_0, d_0 are the entries in Eqn. (5.1.5). Note that $\beta(x) \geq 0$ and it is independent of the earlier choice of unimodular basis. Indeed $\log(a_0 d_0)$ can be expressed geometrically as follows. The decomposition $\mathbf{T}_x Y = W^s(x) \oplus W^u(x)$ and the metric ω_Y determine another metric $\widetilde{\omega}_Y$ defined as the restriction of ω_Y to W^s and W^u and declaring W^s and W^u to be orthogonal. The log of the ratio of volume forms determined by ω_Y and $\widetilde{\omega}_Y$ is exactly $2\beta(x) = \log(a_0 d_0)$. In particular β is also a measurable function, since it is defined using standard constructions on measurable objects.

Finally, we have:

$$(5.1.8) \quad T_Y^* \beta - \beta = \rho^s + \rho^u$$

which says that $\rho^s + \rho^u$ is a coboundary (in the dynamical sense). Indeed, in explicit coordinates as above:

$$\rho^s(x) + \rho^u(x) = \frac{1}{2} \log(a_h d_h) - \frac{1}{2} \log(a_0 d_0)$$

and the formula follows.

5.2. The cohomological equation.

5.2.1. **Proposition** (Cohomologous to a constant). *In the setting of Theorem 1.1.1, there exist $\nu_* \mu$ -measurable functions α_s, α_u such that*

$$(5.2.2) \quad \nu_* \eta_+ |_{\mathcal{W}^u(x)} = e^{-2\alpha_u} \omega_Y |_{\mathcal{W}^u(x)}$$

$$(5.2.3) \quad \nu_* \eta_- |_{\mathcal{W}^s(x)} = e^{-2\alpha_s} \omega_Y |_{\mathcal{W}^s(x)}$$

for $\nu_* \mu$ -a.e. $x \in Y$, where $\mathcal{W}^{s/u}$ denote the stable/unstable manifolds. Moreover the functions satisfy:

$$(5.2.4) \quad \rho^u = T_Y^* \alpha_u - \alpha_u + \frac{h}{2}, \quad \rho^s = T_Y^* \alpha_s - \alpha_s - \frac{h}{2}$$

Proof. To check the existence of functions satisfying Eqn. (5.2.2), Eqn. (5.2.3), it suffices to show that for $\nu_*\mu$ -a.e. x the measure $\nu_*\eta_+|_{\mathcal{W}^u(x)}$ is in the Lebesgue class, since then it must have a Radon-Nikodym derivative $e^{-2\alpha_u}$ relative to the measure $\omega_Y|_{\mathcal{W}^u(x)}$. But by [DTD12, Lemma 5.3] the measures $\nu_*\eta_+|_{\mathcal{W}^u(x)}$ are in the same measure class as the conditional measures of μ along the unstable foliation (regardless of any assumptions on μ). This property on X pushes down to Y . By Ledrappier–Young [LY85, Corollary G], because μ is in the Lebesgue class, the conditional measures along stable manifolds are in the Lebesgue class. The existence of α_u, α_s follows.

The relations in Eqn. (5.2.4) follow from a computation:

$$(5.2.5) \quad e^{2\rho^u - 2T_Y^* \alpha_u} \omega_Y|_{\mathcal{W}^u(x)} = e^{-2T_Y^* \alpha_u} \omega_{Y,h}|_{\mathcal{W}^u(x)} = \\ = \nu_* T^* \eta_+|_{\mathcal{W}^u(x)} = e^h \nu_* \eta_+|_{\mathcal{W}^u(x)} = e^{h-2\alpha_u} \omega_Y|_{\mathcal{W}^u(x)}$$

and similarly for ρ^s . \square

Note also that combining Eqn. (5.1.8) and Proposition 5.2.1 we see that

$$T_Y^*(\alpha_u + \alpha_s - \beta) - (\alpha_u + \alpha_s - \beta) = \rho^u + \rho^s - (\rho^u + \rho^s) = 0,$$

and so by ergodicity of μ :

$$(5.2.6) \quad \alpha_u + \alpha_s = \beta + \delta,$$

μ -a.e., for some constant $\delta \in \mathbb{R}$.

The next proposition is analogous to [CD15, Thm. 5.1], where it is proved by a different argument involving renormalization along the stable manifolds.

5.2.7. Proposition (The restricted current is flat). *For $\nu_*\mu$ -a.e. x we have that*

$$(\xi_x^u)^* \nu_* \eta_+ = e^{-2\alpha_u(x)} \omega_Y(x)|_{W^u(x)}$$

i.e. the pulled back current to the unstable tangent space is a flat metric.

Proof. Consider the function $f(x, r)$ defined for $\nu_*\mu$ -a.e. x and $r > 0$:

$$f(x, r) := \frac{\int_{W^u(x, r)} (\xi_x^u)^* \nu_* \eta_+}{\pi r^2}$$

where $W^u(x, r)$ denotes the ball of radius r for the metric $e^{-2\alpha_u(x)} \omega_Y(x)$ on $W^u(x)$.

By the Lebesgue density theorem and Eqn. (5.2.2), we have that $\lim_{r \rightarrow 0} f(x, r) = 1$ for $\nu_*\mu$ -a.e. x . Applying the automorphism and using the coboundary property, we have that

$$f(T_Y x, e^h r) = f(x, r) \text{ or equivalently } f(T_Y^{-1} x, e^{-h} r) = f(x, r).$$

Now for any $\varepsilon > 0$ there exists $r_\varepsilon > 0$ such that the set

$$Y_\varepsilon := \{x : |1 - f(x, r)| \leq \varepsilon, \forall r \in (0, r_\varepsilon)\}$$

has positive Lebesgue measure. By ergodicity of T_Y , $\nu_*\mu$ -a.e. x visits Y_ε for arbitrarily large negative times, so combined with the above equation it follows that $f(x, r) = 1$ for $\nu_*\mu$ -a.e. x and any $r > 0$.

Take now a point $y \in \mathcal{W}^u(x)$ with $v_y \in W^u(x)$ such that $y = \xi_x^u(v_y)$. We know that the composition $(\xi_y^u)^{-1} \circ \xi_x^u$ is an affine map and let $C_{x,y}$ be the derivative of this map for the metrics $e^{-2\alpha_u(p)}\omega_Y(p)$ with $p = x, y$. It suffices to show that $C_{x,y} = 1$, since then it follows that $\mathcal{W}^u(p)$ carries a canonical flat metric independent of p , and the Radon-Nikodym derivative of $\eta_+|_{\mathcal{W}^u(p)}$ relative to this metric is identically 1 (because $f(x, r) = 1$). By symmetry of the next argument in the two variables, $C_{x,y} \leq 1$ also suffices.

Since $f(y, r) = 1$ for y in a set of full Lebesgue measure on $\mathcal{W}^u(x)$ (and we pick y in this set) it follows that

$$\frac{\int_{W^u(y,r)} (\xi_y^u)^* \nu_* \eta_+}{\pi r^2} = 1 \text{ for all } r > 0.$$

Transporting this identity back to $W^u(x)$ using $(\xi_x^u)^{-1} \circ \xi_y^u$ and using that this map takes a ball of radius r in $W^u(y)$ to a ball of radius $\frac{r}{C_{xy}}$ in $W^u(x)$, it follows that:

$$\frac{\int_{W^u(v_y,r)} (\xi_x^u)^* \nu_* \eta_+}{\pi r^2} = C_{xy}^2$$

where $W^u(v_y, r)$ denotes the ball of radius r at v_y in the metric $e^{-2\alpha_u(x)}\omega_Y(x)|_{W^u(x)}$. But $W^u(v_y, r) \subset W^u(x, r + |v_y|)$ and $\eta_+ \geq 0$ so we have:

$$1 = \frac{\int_{W^u(x,r+|v_y|)} (\xi_x^u)^* \nu_* \eta_+}{\pi(r+|v_y|)^2} \geq \frac{\int_{W^u(v_y,r)} (\xi_x^u)^* \nu_* \eta_+}{\pi(r+|v_y|)^2} = C_{xy}^2 \frac{r^2}{(r+|v_y|)^2}$$

Letting $r \rightarrow \infty$ the desired conclusion follows. \square

5.2.8. Proposition. *The coboundary functions α_u, α_s belong to $L^\infty(\nu_*\mu)$.*

Assuming this, we now finish off the proof of Step 3.

Proof of Theorem 5.1.1. From Proposition 5.2.8 we know that the coboundaries α_u, α_s belong to $L^\infty(\nu_*\mu)$. Since $\rho^s(x) = \log \|DT_Y|_{W^s(x)}\|_{\omega_Y}$, it follows that for all $N \geq 1$ we have

$$\begin{aligned} \log \|D_x T_Y^N|_{W^s(x)}\|_{\omega_Y} &= \rho^s(x) + \cdots + \rho^s(T_Y^{-N+1}x) \\ &= \alpha_s(x) - \alpha_s(T_Y^N x) - \frac{Nh}{2} \\ &\leq -\frac{Nh}{2} + C, \end{aligned}$$

using the coboundary relation (5.2.4) and the L^∞ bound for α_s . This proves (4.1.3), and a similar argument shows (4.1.2).

The L^∞ bound for α_u, α_s together with $\beta = \alpha_u + \alpha_s - \delta$ (by Eqn. (5.2.6)) show that β is uniformly bounded. From Eqn. (5.1.5) and the definition of β in Eqn. (5.1.7), a uniform upper bound on β gives a uniform lower bound on the angle between W^u and W^s , measured relative to ω_Y . \square

5.2.9. Remark. Before the proof of Proposition 5.2.8, note that for μ -a.e. x we have that $(\xi_x^u)^* \nu_* \eta_- = 0$ by say [Can01, Thm. 5.3], or a combination of [DTD12, Cor. 4.2, Prop. 4.4] in the non-projective case. We will restrict

to such x and their image in Y and will eventually construct an entire curve on Y for which the pullback of both $\nu_*\eta_+$ and $\nu_*\eta_-$ vanish. This idea is also a key point in the work of Cantat–Dupont [CD15].

5.2.10. Remark. Recall that the map $\nu: X \rightarrow Y$ contracts the analytic subset $V \subset X$ to the singular points of Y , and that by Proposition 2.1.9 we can write $[\eta_+] + [\eta_-] = \nu^*[\omega_Y]$ for a Ricci-flat orbifold Kähler metric ω_Y on Y . For later use, the following observation will be useful: on Y we can write

$$(5.2.11) \quad \nu_*(\eta_+ + \eta_-) = \omega_Y + i\partial\bar{\partial}\varphi,$$

where $\varphi \in C^0(Y)$.

To see this, we fix smooth representatives α_\pm of $[\eta_\pm]$ on X , and then we can then write $\eta_\pm = \alpha_\pm + i\partial\bar{\partial}\varphi_\pm$ where φ_\pm are Hölder continuous functions on X , as recalled in §2.1.1. On the other hand we can also write $\nu^*\omega_Y = \alpha_+ + \alpha_- + i\partial\bar{\partial}u$ for some continuous function u on X . It follows that

$$\eta_+ + \eta_- = \nu^*\omega_Y + i\partial\bar{\partial}(\varphi_+ + \varphi_- - u),$$

and restricting this to any irreducible component C of V (which as we know is contracted to a point by ν) we have

$$0 \leq (\eta_+ + \eta_-)|_C = i\partial\bar{\partial}(\varphi_+ + \varphi_- - u)|_C,$$

so $\varphi_+ + \varphi_- - u$ is a plurisubharmonic function on the compact curve C , hence constant. It follows that we have $\varphi_+ + \varphi_- - u = \nu^*\varphi$ for some continuous function φ on Y such that $\omega_Y + i\partial\bar{\partial}\varphi$ is a closed positive current on Y which satisfies $\nu^*(\omega_Y + i\partial\bar{\partial}\varphi) = \eta_+ + \eta_-$, and (5.2.11) follows from this.

Proof of Proposition 5.2.8. By Eqn. (5.2.6) together with the fact that $\beta \geq 0$, it suffices to show that $\alpha_u \leq C, \alpha_s \leq C$ on a $\nu_*\mu$ -full measure set. We give the argument for α_u , the one for α_s being identical.

Suppose that α_u is not bounded on the full $\nu_*\mu$ -measure set of points which are not orbifold singularities, for which unstable manifolds exist, and on which Proposition 5.2.7, Proposition 5.2.1, and Remark 5.2.9 hold. So there is a sequence of such points x_i with $\alpha_u(x_i) \rightarrow +\infty$.

Therefore the unstable parametrizations $\xi_{x_i}^u: W^u(x_i) \rightarrow Y$ satisfy $\xi_{x_i}^u(0) = x_i, D_0\xi_{x_i}^u = \mathbf{1}$ and

$$(\xi_{x_i}^u)^*\nu_*\eta_+ = e^{-2\alpha_u(x_i)}\omega_Y(x_i)|_{W^u(x_i)},$$

which goes to zero. Fix a sequence $R_i \rightarrow \infty$.

Suppose first that there is C such that

$$\sup_i \sup_{D_{R_i}(0)} |D\xi_{x_i}^u| \leq C,$$

relative to ω_Y on Y and the flat metric $\omega_Y(x_i)$ on $W^u(x_i)$. Then by Ascoli–Arzelà up to passing to a subsequence, the maps $\xi_{x_i}^u|_{D_{R_i}(0)}$ converge locally uniformly to a nonconstant entire curve $\xi: \mathbb{C} \rightarrow Y$. By construction $\alpha_u(x_i) \rightarrow +\infty$, the current $\nu_*(\eta_+ + \eta_-)$ has locally continuous potentials

(see Remark 5.2.10), and the convergence $\xi_i \rightarrow \xi$ is locally uniform, we can exchange limits to conclude that $\xi^*(\nu_*(\eta_+ + \eta_-)) = 0$.

If there is no such C , we can find points $z_i \in \overline{D_{R_i}(0)}$ such that $|D\xi_{x_i}^u|(z_i) \rightarrow \infty$ (up to subsequence). Working in $D_1(z_i)$, we apply the standard Brody reparametrization argument. Namely for each i we pick a point $y_i \in D_1(z_i)$ which maximizes

$$\delta_i(z) = \text{dist}(z, \partial D_1(z_i)) |D\xi_{x_i}^u|(z).$$

Call $r_i = \text{dist}(y_i, \partial D_1(z_i))$, $a_i = |D\xi_{x_i}^u|(y_i)$, so that

$$a_i \geq a_i r_i \geq \delta_i(z_i) = |D\xi_{x_i}^u|(z_i) \rightarrow \infty,$$

while for all $z \in D_{r_i/2}(y_i)$ we have

$$a_i r_i \geq \text{dist}(z, \partial D_1(z_i)) |D\xi_{x_i}^u|(z) \geq \frac{r_i}{2} |D\xi_{x_i}^u|(z),$$

hence $|D\xi_{x_i}^u| \leq 2a_i$ on $D_{r_i/2}(y_i)$.

Let now

$$\tilde{\xi}_i: D_{a_i r_i/2}(0) \rightarrow Y, \quad \tilde{\xi}_i(z) = \xi_{x_i}^u \left(y_i + \frac{z}{a_i} \right),$$

which are defined on bigger and bigger discs and satisfy

$$\sup_{D_{a_i r_i/2}(0)} |D\tilde{\xi}_i| \leq 2, \quad |D\tilde{\xi}_i|(0) = 1,$$

$$\tilde{\xi}_i^* \nu_*(\eta_+) = \frac{e^{-2\alpha_u(x_i)}}{a_i^2} \omega_Y(x_i)|_{W^u(x_i)},$$

which goes to zero. Again by Ascoli–Arzelà up to passing to a subsequence, the maps $\tilde{\xi}_i$ converge locally uniformly to a nonconstant entire curve $\xi: \mathbb{C} \rightarrow Y$. Since $a_i \rightarrow +\infty$, reasoning as in the case above we again conclude that $\xi^*(\nu_*(\eta_+ + \eta_-)) = 0$.

In both cases, the existence of such an entire curve ξ is a contradiction to Proposition 5.2.12 below. \square

In the proof above we used the following proposition, which is due to Dinh–Sibony [DS05] in general (see [CD15, Proposition 3.10]). In the $K3$ case we can give a very simple proof:

5.2.12. Proposition. *There is no nonconstant entire curve $\xi: \mathbb{C} \rightarrow Y$ such that*

$$\xi^*(\nu_*(\eta_+ + \eta_-)) = 0.$$

Proof. Recall from Remark 5.2.10 that on Y we can write

$$\nu_*(\eta_+ + \eta_-) = \omega_Y + i\partial\bar{\partial}\varphi,$$

where ω_Y is a Ricci-flat orbifold Kähler metric on Y and $\varphi \in C^0(Y)$.

For every $r > 0$ choose χ_r a nonnegative radial cutoff function which equals 1 on D_r , is supported in D_{2r} , and such that $|i\partial\bar{\partial}\chi_r| \leq \frac{C}{r^2}$, for a constant C independent of r . This can be done by letting

$$\chi_r(z) = \eta\left(\frac{|z|}{r}\right),$$

where η is a nonnegative cutoff function on $\mathbb{R}_{\geq 0}$ which equals 1 on $[0, 1]$ and vanishes on $[2, \infty)$. We then compute

$$i\partial\bar{\partial}\chi_r = i\partial\left(\eta' \cdot \frac{z}{2|z|r} d\bar{z}\right) = \left(\frac{\eta''}{4r^2} + \frac{\eta'}{4r|z|}\right) idz \wedge d\bar{z},$$

which is nonzero only for $r \leq |z| \leq 2r$ and so satisfies $|i\partial\bar{\partial}\chi_r| \leq \frac{C}{r^2}$ everywhere. Then for all $r > 0$ we have

$$0 = \int_{D_{2r}} \chi_r \xi^* (\nu_*(\eta_+ + \eta_-)) = \int_{D_{2r}} \chi_r \xi^* (\omega_Y + i\partial\bar{\partial}\varphi),$$

and so the area of $\xi(D_r)$ is bounded above by

$$\int_{D_{2r}} \chi_r \xi^* \omega_Y = \int_{D_{2r}} \chi_r \xi^* (-i\partial\bar{\partial}\varphi) = \int_{D_{2r}} (-\varphi \circ \xi) i\partial\bar{\partial}\chi_r \leq \frac{Cr^2}{r^2} \leq C,$$

and so the entire curve ξ has finite area, hence it extends to a holomorphic map $\xi: \mathbb{P}^1 \rightarrow Y$, see Moncet [Mon12, Prop. 3.31]. If ξ is nonconstant, then by construction the rational curve $C = \xi(\mathbb{P}^1)$ satisfies $C \cdot ([\nu_*(\eta_+ + \eta_-)]) = 0$, which is impossible since $[\nu_*(\eta_+ + \eta_-)] = [\omega_Y]$ is Kähler. \square

Lastly, we prove Corollary 1.1.2:

Proof of Corollary 1.1.2. Cantat has proved [Can99b] that if a compact complex surface X admits an automorphism T with positive topological entropy then X is either a torus, K3, Enriques, a blowup of these, or rational. These are all projective, except for non-projective tori and K3 and their blowups, so thanks to Cantat-Dupont [CD15] (which assume projectivity) we may assume that X is a torus, a K3 surface or a blowup of these.

First, if X is a torus then T is induced by an affine transformation of \mathbb{C}^2 , and therefore (X, T) is trivially a Kummer example. Second, if X is K3 then the result follows from our main Theorem 1.1.1. Lastly, if X is a blowup of a torus or K3 surface, say $\pi: X \rightarrow Y$ is the sequence of blowups, then T induces a bimeromorphic map T_Y of Y , which must be a biholomorphism (see e.g. [BHPV04, Prop. III.4.6]). Furthermore the topological entropy of T_Y equals the one of T by the same argument as in [Kaw08, Lemma 2.2]. From the relation $T_Y \circ \pi = \pi \circ T$ we deduce that the eigencurrents η_{\pm} for T on X are equal to the pullbacks of the corresponding eigencurrents for T_Y on Y , hence the measures of maximal entropy satisfy $\mu = \pi^* \mu_Y$. Since by assumption μ is absolutely continuous with respect to the Lebesgue measure, we conclude that μ_Y is also absolutely continuous with respect to the Lebesgue measure. By the previous cases, we see that (Y, T_Y) is a Kummer example, and hence by definition so is (X, T) . \square

6. ALTERNATIVE ARGUMENTS FOR THE EXISTENCE OF THE COBOUNDARIES

In this section we give an alternative argument for the existence of the expansion/contraction coboundaries in [Proposition 5.2.1](#). Instead of using the theories of Pesin and Ledrappier-Young, we rather exploit our specific geometric setup, and in particular the hyperbolic geometry of the space of Hermitian metrics on a tangent space with given volume form. The coboundaries thus constructed have better integrability properties than those given in [Proposition 5.2.1](#), which are just measurable. We hope that the ideas below may prove useful in related problems.

To simplify notation, in this section we work on X as if it was the orbifold Y , so that $[\eta_+] + [\eta_-]$ is a Kähler class and we will write ω_{Nh} for the Ricci-flat metrics on X that play the role of $\omega_{Y,Nh}$, so that ω_0 replaces ω_Y . In general one would apply the arguments below to the orbifold Y .

6.1. A simple Lemma. The following simple lemma is reminiscent of the Gottschalk–Hedlund theorem.

6.1.1. Lemma (Finding a coboundary). *Let $T: (X, \mu) \rightarrow (X, \mu)$ be a mixing transformation of a probability measure space. Suppose that for $f \in L^2(\mu)$ there exists $C \geq 0$ such that*

$$(6.1.2) \quad \|f + T^*f + \cdots + (T^n)^*f\|_{L^2(\mu)} \leq C,$$

for all $n \geq 1$. Then there exists $h \in L^2(\mu)$ such that $f = h - T^*h$.

Proof. First, note that [Eqn. \(6.1.2\)](#) implies that $\int_X f d\mu = 0$. By the mixing property of T , this implies that the only possible weak limit of $(T^n)^*f$ in $L^2(\mu)$ is 0.

Thanks again to the uniform L^2 boundedness of the Birkhoff sums of f in [Eqn. \(6.1.2\)](#), there is some weak limit $h \in L^2(\mu)$ of $f + T^*f + \cdots + (T^{n_j})^*f$ along some subsequence $\{n_j\}$.

Then using the above remark that the weak limit of $(T^{n_j+1})^*f$ is 0, it follows that

$$(T^*h) = \lim_{n_j} [(f + T^*f \cdots + (T^{n_j+1})^*f) - f] = h - f$$

which is the desired conclusion. □

Recall that automorphisms of $K3$ surfaces with positive entropy are mixing by [[Can01](#), Theorem 6.1].

6.2. Yau vs Oseledets curve.

6.2.1. Hyperbolic geometry. Let V be a complex 2-dimensional vector space, equipped with a non-degenerate complex volume form $\Omega \in \Lambda^2(V^\vee)$. Let $\mathbb{H}^3(V)$ be the space of hermitian metrics which induce the same (real) volume on V as Ω ; this space is naturally isomorphic to real hyperbolic 3-space, since it can be described as the quotient $\mathrm{SL}_2(\mathbb{C})/\mathrm{SU}(2)$.

Consider now a decomposition $V = W_+ \oplus W_-$ into two complex lines. This determines a subset $\gamma \subset \mathbb{H}^3(V)$ of hermitian metrics for which the decomposition of V is orthogonal. There is a natural nearest point projection $\pi_\gamma : \mathbb{H}^3(V) \rightarrow \gamma$, which we now make explicit. Note that γ is a geodesic for the hyperbolic metric.

6.2.2. *Working in coordinates.* Assume that we have $\mathbb{C}^2 = \mathbb{C} \oplus \mathbb{C}$ as our decomposition. A hermitian metric, inducing the standard volume form, is given by a 2×2 matrix:

$$\begin{bmatrix} a & b \\ \bar{b} & d \end{bmatrix} \text{ with } a, d \in \mathbb{R}, b \in \mathbb{C} \text{ and } ad - |b|^2 = 1.$$

The metrics on γ , for which the decomposition is orthogonal, have $b = 0$.

By symmetry considerations, we must have

$$\pi_\gamma \left(\begin{bmatrix} a & b \\ \bar{b} & a \end{bmatrix} \right) = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$

since the map π_γ is equivariant for the action of the matrices in $\text{GL}_2 \mathbb{C}$ which preserve the decomposition $\mathbb{C}^2 = \mathbb{C} \oplus \mathbb{C}$, and the transformation exchanging the two axes is in there.

Finally, using equivariance under the action of scaling the coordinates by $e^{t/2}$ and $e^{-t/2}$ respectively, it follows that:

$$\pi_\gamma \left(\begin{bmatrix} e^{t/2} & 0 \\ 0 & e^{-t/2} \end{bmatrix} \begin{bmatrix} a & b \\ \bar{b} & a \end{bmatrix} \begin{bmatrix} e^{t/2} & 0 \\ 0 & e^{-t/2} \end{bmatrix} \right) = \pi_\gamma \left(\begin{bmatrix} e^t a & b \\ \bar{b} & e^{-t} a \end{bmatrix} \right) = \begin{bmatrix} e^t & 0 \\ 0 & e^{-t} \end{bmatrix}$$

which can be rewritten as

$$\pi_\gamma \left(\begin{bmatrix} a & b \\ \bar{b} & d \end{bmatrix} \right) = \begin{bmatrix} \sqrt{a/d} & 0 \\ 0 & \sqrt{d/a} \end{bmatrix}$$

6.2.3. *Hyperbolic Distances.* Given two Hermitian forms ω_1, ω_2 on V (compatible with the volume) the distance in hyperbolic space between them is defined by

$$\frac{\omega_1 \wedge \omega_2}{\Omega \wedge \bar{\Omega}} = \frac{1}{2} \left(e^{\text{dist}(\omega_1, \omega_2)} + e^{-\text{dist}(\omega_1, \omega_2)} \right).$$

Equivalently, one can pick a basis in which ω_1 is standard Euclidean and diagonalize ω_2 using the spectral theorem to define the distance as the logarithm of the largest (relative) singular value.

6.2.4. *Yau and Oseledets curves.* We now apply the discussion above to $V = \mathbf{T}_x X$ with x in the μ -full measure set where the Oseledets theorem gives us the decomposition $\mathbf{T}_x X = W^s(x) \oplus W^u(x)$. This determines the Oseledets curve $\gamma \subset \mathbb{H}^3(V)$, a hyperbolic geodesic. On the other hand, by Yau's Theorem [Yau78] we have Ricci-flat metrics ω_t on X in the class $e^t[\eta_+] + e^{-t}[\eta_-]$, which together give us the Yau curve in $\mathbb{H}^3(V)$ (although we will be mostly interested only in the values $t = Nh$, $N \geq 1$).

Recall that the expansion/contraction factors defined in Eqn. (5.1.3) are equal to

$$\rho^u(x) = \frac{1}{2} \log \left(\frac{a_h}{a_0} \right), \quad \rho^s(x) = \frac{1}{2} \log \left(\frac{d_h}{d_0} \right)$$

and so a telescoping sum gives

$$\rho^u(x) + \rho^u(Tx) + \cdots + \rho^u(T^{N-1}x) = \frac{1}{2} \log \left(\frac{a_{Nh}}{a_0} \right)$$

and similarly for $\rho^s(x)$.

We will later need the following:

6.2.5. Proposition. *We have the identities:*

$$(6.2.6) \quad \begin{aligned} \int_X \rho^u(x) d\mu(x) &= \frac{h}{2} \\ \int_X \rho^s(x) d\mu(x) &= -\frac{h}{2} \end{aligned}$$

Proof. Recall that the Lyapunov exponent of μ is $h/2$, see §2.2.3.

Once the stable and unstable bundles are given, and since in our case they are line bundles, the exponent can be computed from the formula:

$$\lambda = \int_X \log \|DT|_{W^u(x)}\| d\mu(x)$$

where we compute the norm of $DT: W^u(x) \rightarrow W^u(Tx)$ for one fixed ambient metric. Note that if the metric is changed, then the quantity $\log \|DT|_{W^u(x)}\|$ changes by a coboundary, i.e. $\alpha(x) - \alpha(Tx)$ where $e^{\alpha(x)}$ is the constant of proportionality between the old and the new metric, when restricted to $W^u(x)$. In particular, the integral is independent of the metric.

Since by definition $\rho^u(x)$ is the pointwise norm of DT on the unstable for the Ricci-flat metric, the claimed identity follows. \square

By the discussion in 6.2.1 there are also the ‘‘Oseledets-projected’’ metrics θ_0, θ_{Nh} which correspond to the projection of ω_0, ω_{Nh} to the geodesic γ determined by the stable/unstable decomposition. In the fixed basis $\{e_1, e_2\}$ as in 5.1.4, the metrics are:

$$(6.2.7) \quad \theta_0 = \begin{bmatrix} \sqrt{\frac{a_0}{d_0}} & 0 \\ 0 & \sqrt{\frac{d_0}{a_0}} \end{bmatrix}, \quad \theta_{Nh} = \begin{bmatrix} \sqrt{\frac{a_{Nh}}{d_{Nh}}} & 0 \\ 0 & \sqrt{\frac{d_{Nh}}{a_{Nh}}} \end{bmatrix}$$

The distance-decreasing property of projections in the hyperbolic metric on $\mathbb{H}^3(V)$ gives

$$(6.2.8) \quad \text{dist}(\theta_0, \theta_{Nh}) \leq \text{dist}(\omega_0, \omega_{Nh}) = 2\lambda(x, N).$$

The distance $\text{dist}(\theta_0, \theta_{Nh})$ is computed explicitly as:

$$\begin{aligned}
(6.2.9) \quad \text{dist}(\theta_0, \theta_{Nh}) &= \frac{1}{2} \left| \log \left(\frac{a_{Nh} d_0}{a_0 d_{Nh}} \right) \right| \\
&= \left| \rho^u(x) + \dots + \rho^u(T^{N-1}x) - \rho^s(x) - \dots - \rho^s(T^{N-1}x) \right| \\
&= |S_N \rho^u(x) - S_N \rho^s(x)|
\end{aligned}$$

at μ -a.e. point, where $S_N f$ denotes the Birkhoff sum of the function f ,

$$S_N f(x) := f(x) + \dots + f(T^{N-1}x).$$

Combining this identity with the previous inequality gives

$$(6.2.10) \quad |S_N \rho^u(x) - S_N \rho^s(x)| \leq 2\lambda(x, N)$$

The following observation is going to be crucial:

6.2.11. Proposition. *Suppose that μ is absolutely continuous with respect to the Lebesgue measure, so that $\mu = \frac{d\text{Vol}}{|S|} \Big|_S$. Then the Birkhoff sums of $\rho^u - \rho^s - h$ satisfy*

$$(6.2.12) \quad \int_X e^{S_N(\rho^u - \rho^s - h)} d\mu \leq \frac{2}{|S|}, \text{ for all } N \geq 1.$$

Proof. We have

$$\begin{aligned}
2e^{Nh} &\geq e^{Nh} + e^{-Nh} = \int_X \omega_0 \wedge \omega_{Nh} \\
&= \int_X (e^{2\lambda(x, N)} + e^{-2\lambda(x, N)}) d\text{Vol} \geq \int_X e^{2\lambda(x, N)} d\text{Vol},
\end{aligned}$$

i.e.

$$(6.2.13) \quad \int_X e^{2\lambda(x, N) - Nh} d\text{Vol} \leq 2,$$

but from Eqn. (6.2.10) we also have $S_N \rho^u(x) - S_N \rho^s(x) \leq 2\lambda(x, N)$, and so

$$\int_X e^{S_N(\rho^u - \rho^s - h)} d\text{Vol} = \int_X e^{S_N \rho^u - S_N \rho^s - Nh} d\text{Vol} \leq 2.$$

This finally gives:

$$\int_X e^{S_N(\rho^u - \rho^s - h)} d\mu \leq \frac{2}{|S|}.$$

□

Next, observe that both ρ^u and ρ^s are in L^∞ , with a uniform bound which only depends on the transformation T and the Ricci-flat metric ω_0 . In particular the function $\rho^u - \rho^s - h$ is also in L^∞ , and thanks to Proposition 6.2.5 it satisfies

$$(6.2.14) \quad \int_X (\rho^u - \rho^s - h) d\mu = 0.$$

6.3. Exponential integrability of Birkhoff sums. Our goal now is to use Eqn. (6.2.12) and Eqn. (6.2.14) to prove an exponential integrability bound for $|S_N(\rho^u - \rho^s - h)|$. As it turns out this is an essentially formal consequence of Eqn. (6.2.12) and Eqn. (6.2.14), as we now show.

6.3.1. Proposition (Exponential integrability of Birkhoff sums). *Let $T: (X, \mu) \rightarrow (X, \mu)$ be an invertible transformation of a probability measure space. Suppose that $f \in L^1(X, \mu)$ with $\int_X f d\mu = 0$ is such that*

$$(6.3.2) \quad \int_X e^{S_N f} d\mu \leq C$$

for a uniform constant C and all $N \geq 1$. Then for every $0 < \gamma < \frac{1}{6}$ there is $C' = C'(C, \gamma)$ such that

$$(6.3.3) \quad \int_X e^{\gamma |S_N f|} d\mu \leq C', \text{ for all } N \geq 1.$$

Proof. Decompose $f = f^+ - f^-$ into its positive and negative parts, and for Birkhoff sums denote by $S_N f = S_N^+ f - S_N^- f$ the decomposition into positive and negative parts. Suppose we show that there is a constant C' that depends only on C such that the negative part f^- satisfies the bound

$$(6.3.4) \quad \mu(\{x : f^-(x) \geq L\}) \leq C' e^{-\frac{L}{6}}$$

for all $L \geq 0$.

Then this can be applied to the function $S_N f$ with the transformation T^N , which would thus give us

$$\mu(\{x : S_N^- f(x) \geq L\}) \leq C' e^{-\frac{L}{6}}$$

for all $L \geq 0$ and all $N \geq 1$. On the other hand Eqn. (6.3.2) together with Chebyshev implies that

$$(6.3.5) \quad \mu(\{x : S_N^+ f(x) \geq L\}) \leq C e^{-L}$$

for all $L \geq 0$ and all $N \geq 1$, and so Eqn. (6.3.3) follows from these bounds together with the elementary formula

$$\int_X e^{\gamma |u|} d\mu = \gamma \int_0^\infty \mu(\{|u| \geq t\}) e^{\gamma t} dt.$$

So it suffices to prove Eqn. (6.3.4). Define the set of interest as

$$B_L := \{x : f^-(x) \geq L\}$$

Consider now the set where the positive parts of the Birkhoff sums are large (but on a smaller scale):

$$P_{j,L} := \left\{ x : S_j^+ f(x) \geq \frac{1}{2}(L-1) \right\}$$

which by Eqn. (6.3.5) satisfy $\mu(P_{j,L}) \leq C e^{-\frac{L}{2}}$. Set

$$C_L := B_L \setminus \bigcup_{i=-e^{\frac{L}{6}}}^{e^{\frac{L}{6}}} \bigcup_{j=1}^{e^{\frac{L}{6}}} T^i(P_{j,L})$$

which satisfies the size bound:

$$\mu(C_L) \geq \mu(B_L) - 3e^{\frac{L}{3}} \cdot C \cdot e^{-\frac{L}{2}} = \mu(B_L) - C' e^{-\frac{L}{6}}.$$

Any point $x \in C_L$ has the property that $f^-(x) \geq L$ and for any other point in its orbit $T^i x$ (with $i = -e^{\frac{L}{6}} \dots e^{\frac{L}{6}}$) any Birkhoff sum (with $j = 1 \dots e^{\frac{L}{6}}$) satisfies $S_j^+ f(T^i x) \leq \frac{L-1}{2}$.

Consider now any point y in the support of the function

$$g_L := \mathbf{1}_{C_L} + \mathbf{1}_{T^{-1}C_L} + \dots + \mathbf{1}_{T^{-e^{\frac{L}{6}}}C_L}$$

Then $g_L(y)$ is the number of visits of y to C_L in the times $1 \dots e^{\frac{L}{6}}$. The Birkhoff sum $S_{e^{\frac{L}{6}}}^- f(y)$ can be divided into at most $g_L(y) + 1$ intervals where the positive part is bounded above by $\frac{L-1}{2}$ by the construction of C_L , and the $g_L(y)$ points where $f^- \geq L$. This implies that

$$S_{e^{\frac{L}{6}}}^- f(y) \geq g_L(y) \left(L - 2 \frac{L-1}{2} \right) = g_L(y)$$

for all y in the support of g_L . Integrating over all y gives:

$$\int_{\text{Spt}(g_L)} S_{e^{\frac{L}{6}}}^- f(y) d\mu(y) \geq e^{\frac{L}{6}} \cdot \mu(C_L)$$

On the other hand, Eqn. (6.3.2) implies that $\int_X S_N^+ f d\mu \leq C'$ for a uniform constant C' , but $\int_X f d\mu = 0$ implies $\int_X S_N^- f d\mu = \int_X S_N^+ f d\mu \leq C'$, for all $N \geq 1$. This then implies $\mu(C_L) \leq C' e^{-\frac{L}{6}}$. Using now the lower bound on $\mu(C_L)$ in terms of B_L gives the desired $\mu(B_L) \leq C'' e^{-\frac{L}{6}}$. \square

Thanks to Eqn. (6.2.14) and Eqn. (6.2.12), Proposition 6.3.1 applies to $f = \rho^u - \rho^s - h$, and so we conclude that for every $0 < \gamma < \frac{1}{6}$ there is $C > 0$ such that

$$(6.3.6) \quad \int_X e^{\gamma |S_N(\rho^u - \rho^s - h)|} d\mu \leq C,$$

for all $N \geq 1$. In particular, the hypotheses of Lemma 6.1.1 are satisfied by $f = \rho^u - \rho^s - h$, and we thus obtain

6.3.7. Corollary. *There is a function $\alpha \in L^2(\mu)$ such that*

$$\rho^u - \rho^s - h = T^* \alpha - \alpha.$$

Hence, combining this with Eqn. (5.1.8), the functions

$$\alpha_u := \frac{\alpha + \beta}{2}, \quad \alpha_s := \frac{\beta - \alpha}{2},$$

satisfy Eqn. (5.2.4).

6.3.8. Remark. By working in a suitable Orlicz space instead of $L^2(\mu)$, and applying the analog of Lemma 6.1.1, it is not hard to see that the coboundary α is in fact exponentially integrable, in the sense that $e^{\gamma|\alpha|} \in L^1(\mu)$ for some $\gamma > 0$. With more work, one can deduce the same integrability for β , and hence for α_u and α_s . However, to show that α_u and α_s are in fact bounded, still requires the arguments that we used in Section 5.

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